

Variational Quantum Monte Carlo

① Time Independent Shrodinger Eq.

$$E(\psi) = \int d\vec{R} \underbrace{P(\psi, \vec{R})}_{\substack{\text{Von-Neumann} \\ \text{weights}}} \underbrace{E_L(\psi, \vec{R})}_{\substack{\text{The coordinates of particle} \\ \text{wavefunction}}}$$

Trial/guide wavefunction

$$E(\psi) = \frac{1}{N} \sum_{i=1}^N E_L(\psi, \vec{R}_i)$$

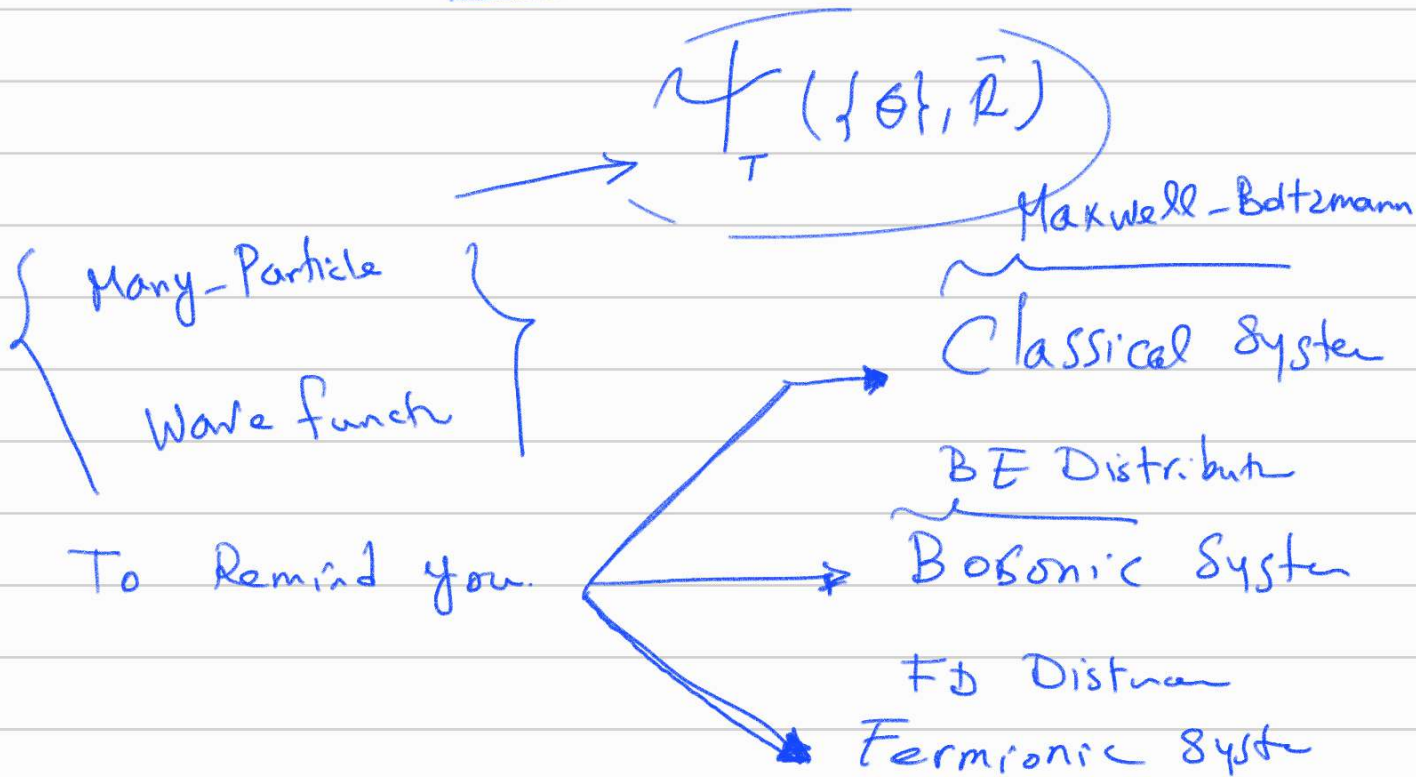
Some Comments on Computing Local Energy


$$E_L(\psi, \vec{R}) = \frac{\mathcal{H} \psi_T(\psi, \vec{R})}{\psi_T(\psi, \vec{R})}$$

$$H = H_{\text{Kinetic}} + H_{\text{Interaction}} = \sum_{i=1}^N H_K^{(i)} + \sum_{i,j} H_{\text{Int}}^{(ij)}$$

$$E_L(\{\theta\}, R) = E_L^{\text{Kinetic}} + E_L^{\text{Interact}}$$

$$\rightarrow E_L^{\text{Kinetic}} = \frac{H_{\text{Kinetic}} \psi_T(\{\theta\}, \vec{R})}{\psi_T(\{\theta\}, \vec{R})}$$







$$\psi_N^{\text{class.}}(\vec{R}) = \prod_{i=1}^N \psi_i^{(i)}(\vec{R}) \rightarrow \text{Many-Body Wave function for Distinguishable}$$

$i=1, 2, 3, \dots, N$

Part 1.

 $\psi_N^{BE}(R) = \psi_N^{\text{Symmetric}}(R) = \sum_p \hat{p} \psi_N^{\text{class}}$

$N=2$ $\psi_2^{BE} = \frac{[\psi_1^{\text{class}}(R_1) \psi_2^{\text{class}}(R_2) + \psi_1^{\text{class}}(R_2) \psi_2^{\text{class}}(R_1)]}{\sqrt{2}}$

 $\psi_N^{FD}(R) = \sum_p \delta_p \hat{p} \psi_N^{\text{class}}$

$$\delta_p = \begin{cases} +1 & \text{even Permutation} \\ -1 & \text{odd Permutation} \end{cases}$$

$\psi_N^{FD}(\bar{R}) \equiv$ Slater determinant

$N=2$ $\psi_2^{FD}(\bar{R}) \equiv \begin{vmatrix} \psi_1(R_1) & \psi_2(R_1) \\ \psi_1(R_2) & \psi_2(R_2) \end{vmatrix}$

$= \frac{\psi_1(\bar{R}_1) \psi_2(R_2) - \psi_1(R_2) \psi_2(\bar{R}_1)}{\sqrt{2}}$

$$E_L^{\text{Kinetic}}(\{t, \vec{R}\}) = ?$$

$$E_L^{\text{Kinetic}}(\{t, \vec{R}_i\}) = \frac{\mathcal{H}_{\text{Kin}}^{(i)} \psi_T(\{t, \vec{R}\})}{\psi_T(\{t, \vec{R}\})}$$

↑
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$$= -\frac{\hbar^2}{2m} \frac{\nabla_i^2 \psi_T(\{t, \vec{R}\})}{\psi_T(\{t, \vec{R}\})}$$

↑
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$$\rightarrow E_L^{\text{Kinetic}}(i) = \underbrace{2T_i}_{\text{Positive}} - \underbrace{|F_i|^2}_{\text{Positive}}$$

$$\rightarrow T_i \equiv -\frac{\hbar^2}{4m} \nabla_i^2 \ln |\psi_T(\{t, \vec{R}\})|$$

$$\rightarrow F_i \equiv \frac{\hbar}{\sqrt{2m}} \vec{\nabla}_i \ln |\psi_T(\{t, \vec{R}\})|$$

$$\lim_{N \rightarrow \infty} (|F_i|^2) = T_i$$

↓
of Run

" if we use above parts to compare "

Local Energy. It can be shown

that, this for may have better

behavior to converge to

expected local Energy

"

"

{ Variational Quantum MC will
Never reach the true ground state }

$$\psi_T(\{\theta\}, \vec{R})$$

A typical form for $\psi_T(\{\theta\}, \vec{R})$

for an Interacting system reads as

Jastrow Correlation factor

$$\Psi_T(\{\theta\}, \vec{R}) = \underbrace{\phi(\{\theta\}, \vec{R})}^{-U(\vec{R})} e^{-U(\vec{R})}$$

Interaction with External field

for Non-Interacting System

$$U(\vec{R}) = \sum_{i=1}^N u_1(\vec{R}_i) + \sum'_{ij} u_2(R_i, R_j)$$

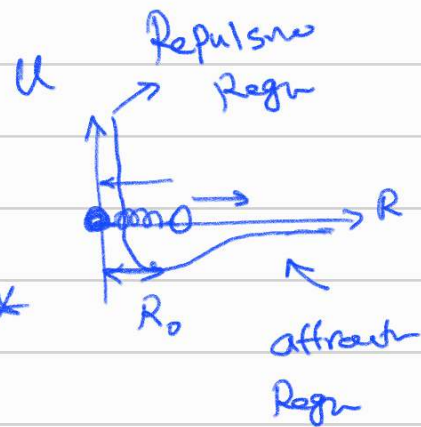
$$+ \sum_{ijk} u_3(R_i, R_j, R_k)$$

$$u_2(\vec{R}_i, \vec{R}_j) = \frac{\theta_1}{R^{12}} - \frac{\theta_2}{R^6} \quad R = |\vec{R}_i - \vec{R}_j|$$

typical Semi-empirical formula

(Lennard-Jones 1924)

$$U(r) = 4\epsilon \left[\left(\frac{\sigma}{r}\right)^{12} - \left(\frac{\sigma}{r}\right)^6 \right]$$



$\{\epsilon, \sigma\} \rightarrow$ it can be determined

via Experiment

$$\mathcal{H} = \sum_{i=1}^N \mathcal{H}_i \Rightarrow E_L = \sum_{i=1}^N E_L^{(i)} = \sum \frac{\mathcal{H}_i \Psi_T}{\Psi_T}$$

$N=2$

$$H_i = \frac{\partial^2}{\partial x_i^2}$$

$$\psi_T = \psi_1 \psi_2 = e^{-\theta_1 x_1} e^{-\theta_2 x_2}$$

$$\frac{H_1 \psi_T}{\psi_T} = \frac{e^{-\theta_2 x_2} \left(\frac{\partial^2}{\partial x_1^2} \right) e^{-\theta_1 x_1}}{e^{-\theta_1 x_1} e^{-\theta_2 x_2}}$$

$$= E_L^{(1)}$$

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$$E_L^{(2)}$$

$$E_L^{(N)} = E_L^{(1)} + E_L^{(2)} + \dots$$

Ex: $N=1$, $D=1$

Simple

harmonic

oscillator

$$H = -\frac{1}{2} \frac{\partial^2}{\partial x^2} + \frac{1}{2} x^2$$

$$\psi_T(\theta, x) = e^{-\theta x^2}$$

$$E(\theta) = \frac{1}{N} \sum_{i=1}^N E_L(\theta, x_i)$$

$$p(\theta, x_i) = \frac{|\Psi_T(\theta, x_i)|^2}{\int dx |\Psi_T(\theta, x)|^2}$$

Subroutine (A)

loop p $\theta = \theta_{min}, \theta_{max}, \Delta\theta$

تکرار متوالی

Loop i=1, N ←

Subroutine (B)

Select x_i from $p(\theta, x_i)$

$$E(\theta) = E(\theta) + E_L(\theta, x_i)$$

End loop

$$E(\theta) = \frac{E(\theta)}{N}$$

Write $\theta, E(\theta)$

End loop

Assume that your Quen has e.g. 3-free

Parameter $\{\theta\}, \{\theta_1, \theta_2, \theta_3\}$

$$\Psi_T = e^{-(\theta_1 x^{\theta_2} + \theta_3)}$$

MCMC

$$E = 0$$

$$\{\theta\}_{old} = \checkmark, E_{old} = \checkmark$$

loop $j=1, M$

(A)

$$\{\theta\}_{old} \xrightarrow{\text{Gaussian funet}} \{\theta\}_{New} = \left\{ \begin{array}{l} \theta_1^{old} + \Delta\theta_1 \\ \theta_2^{old} + \Delta\theta_2 \\ \theta_3^{old} + \Delta\theta_3 \end{array} \right\}$$

$$\Delta\theta_1 = a_1 \mathcal{N}(0, \sigma_1)$$

$$\Delta\theta_2 = a_2 \mathcal{N}(0, \sigma_2)$$

$$\Delta\theta_3 = a_3 \mathcal{N}(0, \sigma_3)$$

$$x_{old} = \checkmark$$

Loop $j=1, N$

$$E(j) = E(i) + E_L(\{\theta\}_{New}, x_{old})$$

$$P(x_{old}) = |4_T(\{\theta\}_{New}, x_{old})|^2$$

$$x_{old} \rightarrow x_{New} = x_{old} + \Delta x$$

$$\Delta x = a_4 \mathcal{N}(0, \sigma_4)$$

$$\rightarrow P(x_{New}) = \checkmark$$

معمولاً مقبول است
 اگر $P(x_{New}) > P(x_{old})$ باشد
 مقبول است
 وگرنه رد می شود

$$AR = \min\left\{1, \frac{P(x_{New})}{P(x_{old})}\right\}$$

$$\left. \frac{P(x_{New})}{P(x_{old})} \right\}$$

(B)

End loop

Metropolis

$$E(j) = \frac{E(i)}{N}$$

$$E_{New} = E(j)$$

$$AR = \min\{1, \frac{E_{Old}}{E_{New}}\}$$

$$\frac{E_{Old}}{E_{New}}$$

(A)

y is Call Random Number

if $y \leq \frac{E_{Old}}{E_{New}}$

$$E_{Old} = E_{New}$$

$$\theta_{Old} = \theta_{New}$$

End if

Metropolis

End loop.

