

# \* Variational Monte Carlo

اصول دروس

① The Variational Principle provide a robust method to find a ground state of typical many-body system

② Consider a typical N-particle Many-Particle Hamiltonian as below:

$$H = H_{\text{Kinetic}} + U \rightarrow \left\{ \begin{array}{l} \text{Kinetic part} \\ \text{Interaction} \end{array} \right.$$

$$H_{\text{Kinetic}} = \sum_{i=1}^N \frac{|\vec{p}_i|^2}{2m} = -\frac{\hbar^2}{2m} \sum_{i=1}^N |\vec{\nabla}_i|^2$$

$$U = \sum_{i=0}^N U_{\text{ext}}(\vec{r}_i) + \sum_{i,j=1}^N U_{\text{int}}(r_{ij})$$

Fermion Anti-Symmetric

Pairwise

Well-Symmetric wavefunction

{ Chapter 10, 11, 12 }  
of Particle

Boson fully symmetric

③

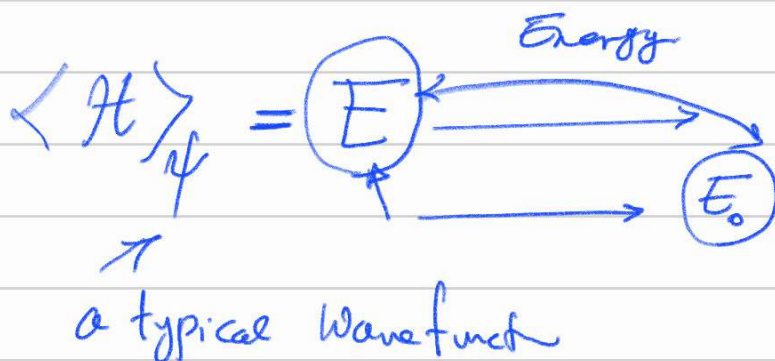
$$\boxed{H \psi_n = E_n \psi_n}$$

$\psi_n$  is the nth of eigen wave function

$E_n$  is eigen value

$$\psi_n?, E_n? \rightarrow \begin{cases} E_0 \\ \psi_0 \end{cases} \text{ Ground State}$$

- ④ The Purpose of variational approach is to find a wave function that can be optimized to produce a closest expected value of  $\langle H \rangle$  to the ground state



We are interested in finding a method

to map the eigen value problem to

importance sampling approach

$E_0$

Variational MC

⑤ Example: Time-Independent Schrödinger Equation

Quantum Harmonic Oscillator in 1D

$$\begin{cases} \hat{H}\psi = \left[ -\frac{\hbar^2}{2m} \frac{d^2}{dx^2} + \frac{1}{2} m \omega^2 x^2 \right] \psi \\ \hat{H}\psi = E\psi \end{cases}$$

$H_0(x) = 1, H_1(x) = 2x \dots$

$E_n = (n + \frac{1}{2}) \hbar \omega$

Hermite Polynomial

$$\psi_n(x) = \left( \frac{m\omega}{\pi\hbar} \right)^{1/4} \frac{1}{\sqrt{2^n n!}} H_n \left( x \sqrt{\frac{m\omega}{\hbar}} \right) e^{-\frac{m\omega x^2}{2\hbar}}$$

⑥ Variational Principle

$\hat{H}\psi_n(x) = E_n \psi_n(x)$

$\begin{cases} E_n = \text{eigenvalue} \\ \psi_n = \text{eigenfunction} \end{cases}$

$\hat{H}\psi_0(x) = E_0 \psi_0(x)$

$\begin{cases} E_0 = ? \\ \psi_0 = ? \end{cases}$

$\psi(x)$  : a typical arbitrary wave function

مثال:  $\psi(x) = e^{-\alpha|x|}$



$$\langle \mathcal{H} \rangle_{\psi} = E$$

bases func

$$\psi(x) = \sum_n C_n \phi_n$$

تقریباً صحیح دیتا

$$E = \langle \mathcal{H} \rangle_{\psi} = \frac{\langle \psi | \mathcal{H} | \psi \rangle}{\langle \psi | \psi \rangle}$$

$$E = \frac{\int dx \psi^*(x) \mathcal{H} \psi(x)}{\int dx \psi^*(x) \psi(x)}$$

In principle, this value can be as close as we want to  $E_0$  of  $\mathcal{H}$

$$E \rightarrow E_0$$

$$\left\{ \begin{array}{l} E = \langle \mathcal{H} \rangle_{\psi} \geq E_0 \quad \text{for any arbitrary } \psi \\ \text{and } \boxed{E = E_0} \quad \text{if and only if} \\ \psi = C_0 \phi_0 \end{array} \right.$$

$$H \varphi_n(x) = E_n \varphi_n(x)$$

$$\int \varphi_n^*(x) \varphi_m(x) dx = \delta_{nm}$$

$$\psi(x) = \sum_n C_n \varphi_n(x)$$

$$E = \langle H \rangle_\psi = \frac{\int dx \psi^*(x) H \psi(x)}{\int dx \psi^*(x) \psi(x)}$$

$$E = \frac{\sum_{nm} C_n^* E_n C_m \int dx \varphi_n^*(x) \varphi_m(x)}{\sum_{nm} C_n^* C_m \int dx \varphi_n^*(x) \varphi_m(x)}$$

$$\sum_{nm} C_n^* C_m \int dx \varphi_n^*(x) \varphi_m(x)$$

$$E = \frac{\sum_{n=0}^{\infty} |C_n|^2 E_n}{\sum_{n=0}^{\infty} |C_n|^2} = E_0 + \frac{\sum_{n=1}^{\infty} |C_n|^2 E_n}{\sum_{n=0}^{\infty} |C_n|^2}$$

$$E = E_0 + \frac{\sum_{n=0}^{\infty} |C_n|^2 (E_n - E_0)}{\sum_{n=0}^{\infty} |C_n|^2}$$

$E_n \gg E_0$   
 $n=0, 1, 2, \dots$

$$E = E_0 + \overbrace{\dots}^{\text{positive definite}}$$

$$E \geq E_0$$

The minimum value of our estimated 'E' is an upper limit of ground state.

$$\textcircled{7} \quad \psi(x; \{\theta\}) \longrightarrow \underbrace{\psi(x; \{\theta\}_{\text{Best}})}_{\text{Min}\{E\} = E_0}$$

$E \geq E_0$

Variational MC

⑧ Variational MC

قدر پارامترها را در تابع قرار می‌دهیم

$$\textcircled{A} \quad \psi_T(\{\theta\}, x) = \psi_T(\{\theta_1, \theta_2, \dots, \theta_n\}, x)$$

(B)

$$E(\psi_t) = \langle \mathcal{H} \rangle_{\psi_T}$$

$$= \frac{\int dx \psi_T^*(\psi_t, x) \mathcal{H} \psi_T(\psi_t, x)}{\int dx \psi_T^*(\psi_t, x) \psi_T(\psi_t, x)}$$

$\{x\} \equiv \{\bar{r}_1, \bar{r}_2, \dots, \bar{r}_N\}$

$N \equiv$  # of particles in system

(C)

$$P(\psi_t, x) \equiv$$

$$\frac{|\psi_T(\psi_t, x)|^2}{\int dx |\psi_T(\psi_t, x)|^2}$$

$$|\psi_T(\psi_t, x)|^2$$

(D)

(B) and (C)  $\rightarrow$

$$E(\psi_t) = \frac{\int dx \psi_T^*(\psi_t, x) \frac{\psi_T(\psi_t, x)}{\psi_T(\psi_t, x)} \mathcal{H} \psi_T(\psi_t, x)}{\int dx |\psi_T(\psi_t, x)|^2}$$


$$= \frac{\int dx |\psi_T(\psi_t, x)|^2 \frac{\mathcal{H} \psi_T(\psi_t, x)}{\psi_T(\psi_t, x)}}{\int dx |\psi_T(\psi_t, x)|^2}$$

$$E(\psi_t) = \int dx P(\psi_t, x)$$

$$\frac{\mathcal{H} \psi_T(\psi_t, x)}{\psi_T(\psi_t, x)}$$



$$E(f(\theta)) = \int dx p(f(\theta), x) E(f(\theta), x)$$

$\int dx A(x)$   
 Probability  


Recall  $\langle f \rangle = \int dx f p(x) = \frac{1}{N} \sum_{i=1}^N f(x_i)$

$x_i$  can be generated with probability  $p(x)$

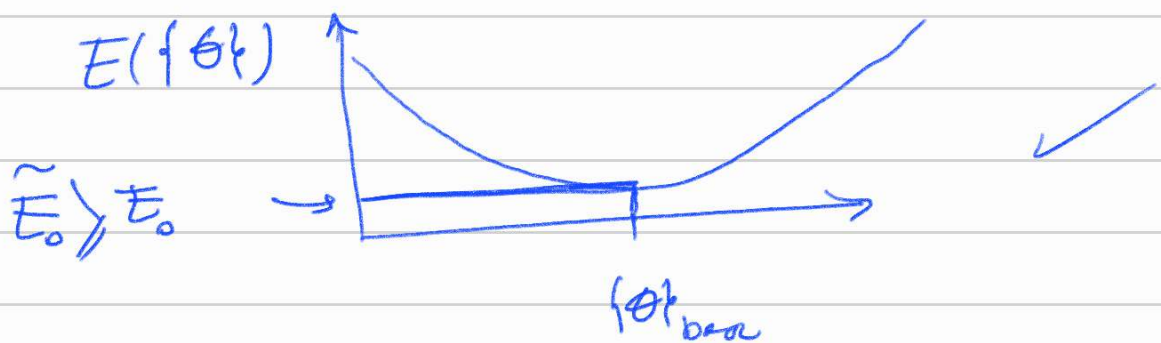
$$E(f(\theta)) = \frac{1}{N} \sum_{i=1}^N E(f(\theta), x_i) p(f(\theta), x_i)$$

$x_i$  با احتمال  $p$  تولید می شود، پس وزن  $p$  می دهد.

$$E(f(\theta)) = \frac{1}{N} \sum_{i=1}^N E(f(\theta), x_i)$$

for a given set of  $f(\theta)$ .

$p(f(\theta), x)$





\* for a given set of  $\{\theta\}$ . \*

We should produce of  $X$  with

Probability given by  $P(\{\theta\}, X)$

\* and Compute  $E(\{\theta\})$  \*

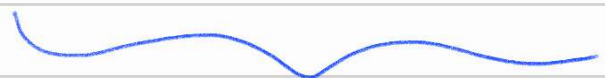
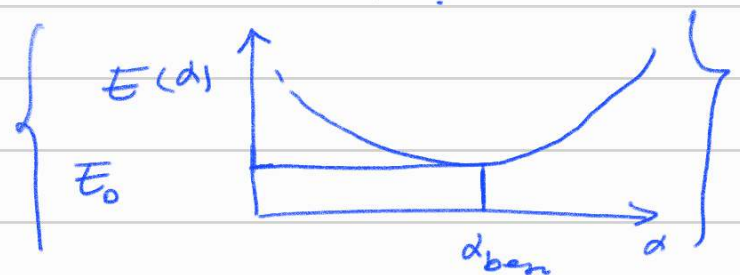
Ex 2  $H = -\frac{1}{2} \frac{d^2}{dx^2} + \frac{1}{2} x^2$   $\left\{ \begin{array}{l} m=1 \\ \hbar=1 \\ \omega=1 \end{array} \right\}$

$$\left\{ \begin{array}{l} \psi_0 = ? \\ E_0 = ? \end{array} \right.$$

We choose the following trial wave function to solve our problem

$$\psi_T(\theta, x) = e^{-\frac{1}{2} \alpha x^2}$$

$\alpha = ?$





End loop

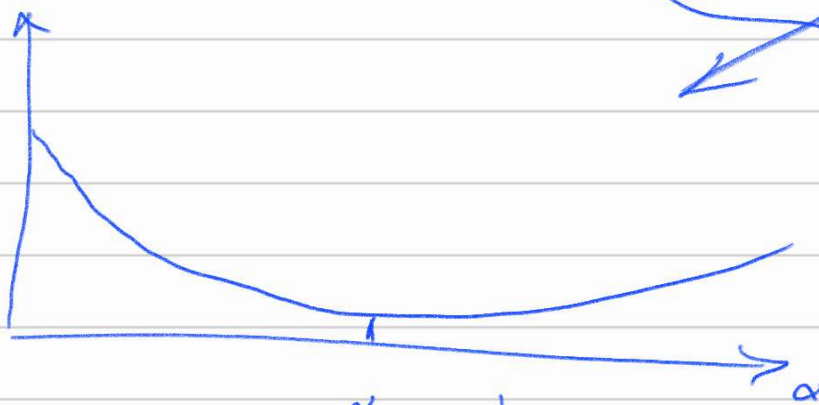
$$E(\alpha) = \frac{E(\alpha)}{N}$$

Write  $\alpha$ ,  $E(\alpha)$

End loop

Atypical VMC algorithm

$E(\alpha)$   
↑



$$\alpha_{\text{best}} = 1/2$$

→ Analytical solution

Importance Sampling + Metropolis

+ MCMC (HMC)

$\{X\}$ ,  $\{\theta\}$